



### **OPTICAL MATERIALS**

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### Vision – our most important tool for perceiving our surroundings

- evolved several hundred million years ago
- vital not only for humans
- high resolution imaging perception
  - Far field resolution: ca. 1 arc minute
  - Near field resolution: ca. 100 μm
- Iong range (millions of light years)

#### Phrases (german and english):

- "Sich ein Bild machen"
- "Ein Bild sagt mehr als tausend Worte"
- "Seeing is believing"







### 1.1 Significance of light Expanding our capabilities of perception



Optical instruments for a better understanding of the world we live in







Telescopes help us to understand our universe



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3-D laser imaging of cell www.leibniz-inm.de









### 1.2 What is light? Comparison of theories

### Particles (classical):

- explains straight line propagation
- explains reflection (elastic bouncing)
- has difficulties with refraction
- does not explain diffraction & related phenomena

### Waves (classical):

- less intuitive
- successfully explains all phenomena of classical optics

#### Quantum mechanics:

- answer depends on how the question is asked
- unambiguously particle-like (photons):

quantization of energy hf and momentum  $\frac{h}{2\pi}k$ 



# 1.3 Wave theory The wave equation (1)

### Wave propagation

= specific mode of evolution of a quantity depending on position and time

("field", "amplitude")

### Examples:

- Height of water surface on the sea
- Pressure (sound = acoustic waves)
- Electric field  $\vec{E}$  and magnetic field  $\vec{H}$  (light)

### Wave equation

Description of wave propagation for a generalized field A:

$$\frac{\partial^2 A}{\partial x^2} + \frac{\partial^2 A}{\partial y^2} + \frac{\partial^2 A}{\partial z^2} = \frac{1}{v^2} \frac{\partial^2 A}{\partial t^2}$$

For electromagnetic waves: a consequence of Maxwell's equations





 1.3 Wave theory The wave equation (2)



Simplest case – only one spatial dimension :

$$\frac{\partial^2 A}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 A}{\partial t^2}$$

~? .

General solutions:

 $A(x,t) = A(x - v \cdot t)$  or  $A(x,t) = A(x + v \cdot t)$ (v = velocity of propagation)



### 1.3 Wave theory Harmonic waves (one-dimensional)



### Harmonic (sine/cosine) dependency on position and time:

Amplitude of a harmonic wave propagating in +x direction:

$$A(x,t) = A_0 \cdot \cos\left[\frac{2\pi}{\lambda}(x-\nu \cdot t)\right] = A_0 \cdot \cos(k \cdot x - \omega \cdot t) = A_0 \cdot \cos(\varphi)$$

Where

 $\lambda = wavelength (distance between two repeat units)$ 

ω = 2πf = <sup>2πν</sup>/<sub>λ</sub> = angular frequency (2π times the oscillation frequency f)
 k = <sup>2π</sup>/<sub>λ</sub> = wave number = number of repeat units on 2π length units
 φ = k · x - ω · t = phase of the harmonic function

### 1.3 Wave theory Generalizations

Plane waves in 3D space:

Wave number  $k \rightarrow$  wave vector  $\vec{k}$ 

- Points along direction of propagation
- Length =  $2\pi$ /wavelength\*
- Wave fronts (planes of constant phase) are perpendicular to  $\dot{k}$

\*true for individual vector components as well

Using complex numbers:  $\exp(i \cdot \varphi) = \cos(\varphi) + i \cdot \sin(\varphi)$  $A(x,t) = Re\left[A_0 \cdot \exp\left\{i \cdot (\vec{k} \cdot \vec{x} - \omega \cdot t)\right\}\right] = Re[A_0 \cdot \exp(i \cdot \varphi)]$ 

Why: Adding phases is simpler, generalized formalism covers more cases





1.4 Light as an electromagnetic wave Electric and magnetic fields



### Electric field $\vec{E}$

• Defined as the force  $\vec{F}$  acting on a test charge q, scaled by the magnitude of that test charge:

$$\vec{F} = q \cdot \vec{E}$$

- The electric field has a strength and a direction  $\rightarrow$  **vector**
- It depends on the location in space (coordinate vector  $\vec{x}$ )  $\rightarrow$  field (and, generally, on time t as well):

 $\vec{E} = \vec{E}(\vec{x},t)$ 

## Magnetic field $\vec{H}$ (or $\vec{B} = \mu \cdot \mu_0 \cdot \vec{H}$ )

- Can be defined similarly (difficulty: there are no "magnetic charges")
- Usually not interesting for optics: magnetic permeability  $\mu$  is 1 for almost all materials at optical frequencies



- Density  $\overrightarrow{P}$  of induced dipoles (dielectric polarization) is proportional to the electric field:  $\overrightarrow{P} = \chi \cdot \overrightarrow{E}$
- $\chi$  (the dielectric susceptibility) is a real number if the material's reaction to the electric field is instantaneous
- $\chi$  has a (frequency dependent) imaginary component if the reaction lags behind ( $\rightarrow$  absorption)
- Important quantity: (relative) dielectric permittivity  $\varepsilon = \chi + 1$

## 1.4 Light as an electromagnetic wave Electromagnetic waves

Wave equation for electric field (from Maxwell's equations):

$$\frac{\partial^2 \vec{E}}{\partial x^2} + \frac{\partial^2 \vec{E}}{\partial y^2} + \frac{\partial^2 \vec{E}}{\partial z^2} = \varepsilon \varepsilon_0 \mu \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2} = \frac{n^2}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}$$

where

- $c = \sqrt{\frac{1}{\varepsilon_0 \cdot \mu_0}}$  (= 299 792 458 m/s) speed of light (in vacuum) ( $\varepsilon_0, \mu_0$  = vacuum permittivity / permeability)
- $n = \sqrt{\varepsilon \cdot \mu}$  = refractive index (of a material) ( $\varepsilon, \mu$  = relative permittivity / permeability for a given material)
- $\rightarrow$  A material slows down light by a factor of n

Harmonic waves:  $\lambda$  and  $\vec{k}$  depend on the material:  $\lambda \rightarrow \frac{\lambda}{n}$ ,  $\vec{k} \rightarrow n \cdot \vec{k}$ 





### 1.4 Light as an electromagnetic wave Harmonic electromagnetic waves in isotropic media

### Relation between wave vector and field vectors

- Electric field  $\vec{E}$  and magnetic field  $\vec{H}$  are perpendicular to each other and to  $\vec{k}$
- Amplitudes are linked:  $\|\vec{E}\| = \sqrt{\frac{\mu\mu_0}{\varepsilon\varepsilon_0}} \cdot \|\vec{H}\|$

#### Polarization:

- Linear polarization:  $\vec{E}$  and  $\vec{H}$  oscillate in orthogonal planes
- Circular / elliptical polarization is also possible

### Intensity:

- Intensity *I* = power per unit area perpendicular to  $\vec{k}$
- $I = \|\vec{E} \times \vec{H}\|$ : proportional to the squared amplitude





Loo Kang Wee (via Wikipedia)

### 1.4 Light as an electromagnetic wave Light in the electromagnetic spectrum





#### Visible light: $\sim 380 - 780$ nm

### 1.4 Light as an electromagnetic wave Interfaces between different materials



### Continuity conditions (required by Maxwell's equations)

- Tangential components of  $\vec{E}$  and  $\vec{H}$  are continuous, i.e., the same on both sides of the interface
- Perpendicular components of  $\vec{D} = \varepsilon \varepsilon_0 \vec{E}$  and  $\vec{B} = \mu \mu_0 \vec{H}$  are continuous

### Consequence (a): Snell's law of refraction

 Continuity conditions are true at every point of the interface

 $\rightarrow$  tangential component of  $\vec{k}$  is continuous Length of  $\vec{k}$  changes according to different refractive index

 $\rightarrow$  perpendicular component must adapt





### 1.4 Light as an electromagnetic wave Interfaces between different materials



### Consequence of continuity conditions (b): Fresnel coefficients

Three waves to match all conditions : incident, transmitted, reflected 

 $\alpha_1$ 

Matching conditions depend on polarization 



plane of incidence is spanned by 3 wave vectors

polarization is defined by  $\vec{E}$  field with respect to the plane of incidence

Amplitude reflection and transmission coefficients: 

$$r_{s} = \frac{n_{1} \cos \alpha_{1} - n_{2} \cos \alpha_{2}}{n_{1} \cos \alpha_{1} + n_{2} \cos \alpha_{2}} \qquad r_{p} = \frac{n_{1} \cos \alpha_{2} - n_{2} \cos \alpha_{1}}{n_{1} \cos \alpha_{2} + n_{2} \cos \alpha_{1}}$$
$$t_{s} = \frac{2n_{1} \cos \alpha_{1}}{n_{1} \cos \alpha_{1} + n_{2} \cos \alpha_{2}} \qquad t_{p} = \frac{2n_{1} \cos \alpha_{1}}{n_{1} \cos \alpha_{2} + n_{2} \cos \alpha_{1}}$$

reflected:

$$E_r = r \cdot E_i$$

transmitted:  $E_t = t \cdot E_i$ 

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### Superposition of waves $\rightarrow$ amplitudes add up

1.4 Light as an electromagnetic wave

#### Constructive interference:

Interference

- same direction / sign
- $I \propto (E_1 + E_2)^2 = E_1^2 + E_2^2 + 2E_1E_2$
- Intensity > sum of separate intensities

#### Destructive interference:

- opposite direction / sign
- $I \propto (E_1 E_2)^2 = E_1^2 + E_2^2 2E_1E_2$
- Intensity < sum of separate intensities



temporal: sufficiently stable frequency for paths of different length

spatial: different points of a light source are correlated





### 1.4 Light as an electromagnetic wave Huygens' principle



### Not everything is a plane wave

Basic idea: Each point in space is the origin of a spherical wave correlated with the wave field. All spherical waves interfere with each other.

- Undisturbed case: The result is the wave field as it was
- Disturbed case: Basis for scalar diffraction theory (Perturbation: any inhomogeneity in the material, i.e. the spatial distribution of ε)
- ► Example: Young's double-slit experiment
   → ultimate proof for wave character of light



## 2.1 Optical effects of periodic microstructures Gratings and lattices



### Description of periodic structures in space

Grating: 2D (sheet-like) structure, periodic within the sheet plane Lattice: 3D (volume) structure, periodic in all three dimensions

Simplest case: grating of parallel lines

Just like an optical wave, but static:

 $F(\boldsymbol{x}) = A \cdot \cos\left(\vec{K} \cdot \vec{x}\right)$ 

 $\vec{K}$  = grating vector (or lattice vector)

*F* can be any quantity relevant to optics:

- refractive index
- absorption
- height of a surface (2D gratings only)



### 2.1 Optical effects of periodic microstructures Grating diffraction

### Interaction of a light wave with a grating

Incident light:

$$\vec{E}_{incident}(\vec{x},t) = \vec{E}_{incident}^{0} cos\left(\vec{k}\cdot\vec{x}-\omega t\right)$$

 $\vec{k}$  is not parallel to the grating plane

### Result:

- Attenuation of incident wave
- Additional plane waves:

$$\vec{E}_{\pm}(x,t) = \vec{E}_{\pm}^{0} \cos\left(\vec{k}_{\pm} \cdot \vec{x} + \delta_{\pm} - \omega t\right)$$

- tangential component of  $\vec{k}_{\pm}$  is obtained by adding / subtracting  $\vec{K}$
- perpendicular component of  $\vec{k}_{\pm}$  adjusts to maintain correct length





2.1 Optical effects of periodic microstructures Diffraction – general case



### General description of a periodic structure:

Superposition of harmonic modulations with different K (Fourier representation)

Condition for diffraction: There must be a lattice vector  $\vec{K}$  present with

$$\vec{k}_{diffracted} = \vec{k}_{incident} \pm \vec{K}$$
 (Bragg condition)

Additionally, the length of  $\vec{k}$  must remain constant (conservation of energy).

- Diffraction only for special combinations of incident and diffracted wave vectors
- Basis for crystal structure analysis by X-ray diffraction

Special case for thin gratings: perpendicular component of  $\vec{k}_{\pm}$  adjusts to maintain correct length

### 2.1 Optical effects of periodic microstructures Examples





hexagonal array of holes in a metal film



CD / DVD



natural opal

2.2 Non-periodic microstructures From diffraction to scattering



Loss of strict periodicity in the structure

→ Fourier representation turns into a continuous spectrum

 $\rightarrow$  Loss of selectivity in the diffraction condition

Partial loss of periodicity:

- Still some selectivity
- Preferences for some wavelengths: structural colors
- Preferences for some angles

Complete loss of periodicity / complete randomness:

- Uniform scattering
- Can be interpreted as a "random walk" of photons



### 2.2 Non-periodic microstructures Examples



#### Partially periodic





#### **Completely random**





### 2.3 Periodic nanostructures Metamaterials



#### Transition from "micro" to "nano":

- Semi-official: feature size < 100 nm
- Convenient for optics: periodicity <  $\lambda/2$  (roughly equivalent for visible light)

#### Why?

- Periodicity <  $\lambda/2$  means  $K > 2 \cdot k$
- No chance to maintain correct length of the diffracted wave vector, even with a full reversal of the direction of propagation
- Result: periodic structure becomes invisible!

But: Properties of the structural units still remain active

- Anisotropy & polarization
- Unusual dispersion relations (wavelength dependence of refractive index)

### 2.3 Periodic nanostructures Metamaterials - Examples



#### Moth eye: one of few examples for optical wavelengths





#### Many proof of principle demonstrations with microwaves

#### Cloaking devices

. . .

Negative refractive index



Metamaterial cloak (Source: K. Kim, Yonsei University; via Welt der Physik)

### 3.1 Vertically structured surfaces Modifying reflectivity

#### Surface reflection from transparent materials

Reason: Matching electromagnetic field amplitudes on both sides of the surface (Fresnel reflection)

Amplitude reflection coefficient:

$$r = \frac{n_1 - n_2}{n_1 + n_2}$$

@ normal incidence (angle dependent) (Intensity reflection coefficient is  $r^2$ !)

### Task: Either decrease or increase reflectivity







3.1 Vertically structured surfaces Why antireflective surfaces?

Maximize light throughput:



#### Maximize contrast:









3.1 Vertically structured surfaces
 Approaches for antireflective surfaces (2)



Gradient index (GRIN) surfaces



True GRIN is difficult, approximation by a series of layers with small refractive index differences is possible.

Additional difficulty: Refractive index of air is 1.0, but the lowest refractive index of a solid material is 1.38 ( $MgF_2$ ); porous silica can go down to 1.22, but has other issues.

#### Not competitive in practice

Layer parameters for optimum antireflective effect:

Destructive interference of reflected partial waves

Approaches for antireflective surfaces (2)

3.1 Vertically structured surfaces

refractive index =  $\sqrt{n_1 n_2}$ 



Single interference layer



Technically used for low cost applications



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3.1 Vertically structured surfaces
 Approaches for antireflective surfaces (3)

### Multiple interference layers

Increased bandwidth

Various designs possible (example: 3 layers) Still, generally: thickness of each layer =  $\lambda/4$ 







3.1 Vertically structured surfaces
 Approaches for antireflective surfaces (4)

#### Moth eye structures

Metamaterial to achieve

- Iow average refractive index
- some GRIN effect



300 nm



Rather limited technical application sensitive to soiling and wear



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### General requirements for layer materials:

- Refractive index: from very low to very high
- Transmission range may vary depending on the application
- Resistance to environmental conditions depending on the application

### Typical materials:

- MgF<sub>2</sub> (n=1.38, soft, broad transmission range)
- Amorphous silica (SiO<sub>2</sub>; n= 1.46, hard, UV transparent)
- Porous silica (n down to 1.22, very soft)
- TiO<sub>2</sub> (n = 2.4 to >3 for dense material; hard; UV cutoff below 400 nm)
- ZrO<sub>2</sub> (n = 2.13; hard; transparent for near UV)

### 3.1 Vertically structured surfaces How to make interference layer systems

### 1. Gas phase deposition

- Physical vapor deposition
- Chemical vapor deposition

#### Features:

- High quality coatings
- Suitable for small substrates
- High equipment cost







### 2. Sol-Gel and related methods

Wet chemical synthesis based on hydrolysis and condensation reactions starting from suitable precursors



- Depending on synthesis route: nanoparticles or amorphous network dispersed in organic solvent
- May be modified with organic cross-linkers for low-temperature stability
- Various coating methods: Spin coating, dip coating, continuous roll-to-roll processes
- UV and / or thermal curing

3.1 Vertically structured surfaces Wet coated interference layers

Dip coating for rigid substrates



Roll-to-roll coating line





Coated PET foil



beam splitter

laser beam

### Origination methods for nano-/microstructured surfaces

Self assembly of nano-/microparticles

mirror

3.2 Horizontally structured surfaces

Gratings, holograms and moth eyes

- Specific etching methods
- ▶ Direct writing (laser, electron beam) → high resolution, but very slow

#### Holographic writing:



substrate

coated with







3.2 Horizontally structured surfaces From master structure to mass production



Origination → Master structure



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### 3.2 Horizontally structured surfaces Replication techniques

- Hot embossing
  - limited to thermoplastic polymers
  - needs to cool while in contact with the tool
    → slow process
- Reactive casting
  - UV curing through the substrate foil (needs to be transparent) or through a transparent tool (silicone)
  - well-established process, fast
- Embossing into a thixotropic resin (INM approach)
  - shaping with high shear rate, but no relaxation afterwards
  - curing after removing the tool
  - fast alternative for non-transparent substrates







Goal: use nanoparticles to modify the properties of a polymer material ...

- Change the refractive index
- Modify the dispersion curve
- Increase hardness / wear resistance
- Add electrical functionalities
- ...

... without sacrificing transparency!

Challenge: Periodic structures may become invisible with feature sizes below  $\lambda/4$  (half period), but inhomogeneities in random distributions of particles are much larger than the particles themselves  $\rightarrow$  Particles need to be really small!

### 4. Nanocomposites Importance of particle size

Transparency of randomly distributed particles:

$$\frac{I}{I_0} = exp -4 \cdot \frac{\pi^4}{\lambda^4} \cdot d^3 \cdot \left(\frac{n_p^2 - n_m^2}{n_p^2 + 2 \cdot n_m^2}\right)^2 \cdot c \cdot L$$



- $\lambda$  = Wavelength
- d = Particle diameter
- $n_p$  = Refractive index particles
- $n_{m}^{'}$  = Refractive index matrix
- c = Particle concentration
- L = Thickness bulk

- Notes: 1. Applies also to nanoporous materials
  - 2. Good dispersion (avoiding agglomeration) is equally important

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### • 4. Nanocomposites

An example of a nanocomposite for microstructures

#### Target

- Development of Light Management Foils (LMF)
- Enhancement of brightness and contrast, reduced viewing angle dependence for LCD
- Better brightness and contrast, lower sensitivity to ambient light for projection screens

### Methods

- Photosensitive gradient index material based on cross-linkable nanoparticles in a gel-like matrix
- Irradiation of this material through a mask produces a columnar microstructure with angledependent scattering properties
- Continuous roll-to-roll processes for coating, mask lamination and irradiation









### 4. Nanocomposites Light management foils

### Results

- 50 μm thick films with pronounced angle-dependent scattering
  - High haze (>94 %) for light incident from preferred direction
  - Significantly lower haze for other directions
- LMF as diffuser in LCDs:
  - approx. 20 % higher brightness and contrast
- LMF on mirror delivers even greater improvement

### Applications

- Diffusers for LCD panels
- Projection screens
- Lighting

Projection screen LMF on mirror





### Summary of most important topics



#### Theory

- What is light
- Formalism for describing electromagnetic waves
- Interaction of light with materials structured on microscopic / submicroscopic scale

### Application & technology

- Antireflective surfaces
- Microstructured surfaces
- Nanocomposites



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